

A New Way to Measure Spin at Hadron Colliders

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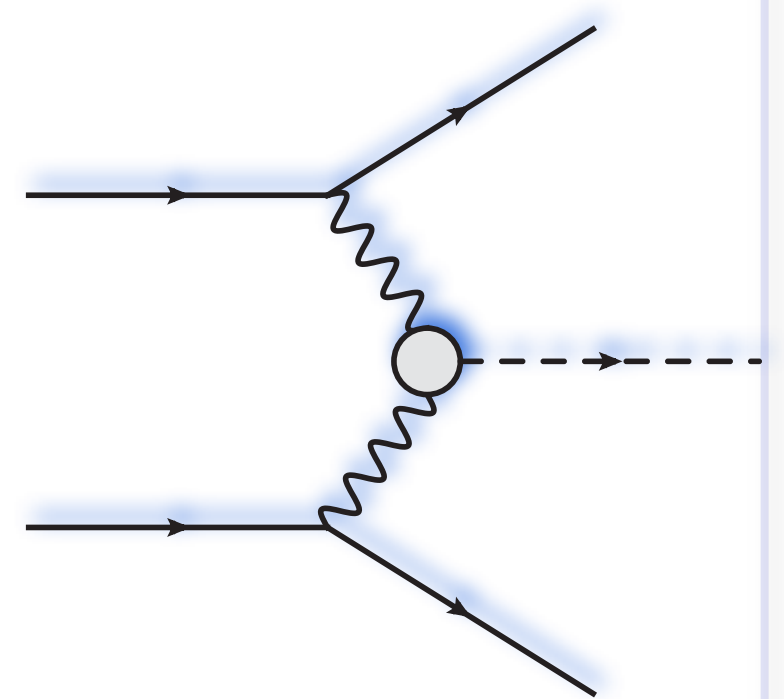
with Michael Ramsey-Musolf 1008:5151

Spin Determination

- Many SM extensions contain new strongly interacting particles that decay into SM + missing energy
- Spin measurements key to distinguishing possibilities
- Ideally, want a technique that doesn't rely on long decay chains, chiral couplings, or decays into specific final states.

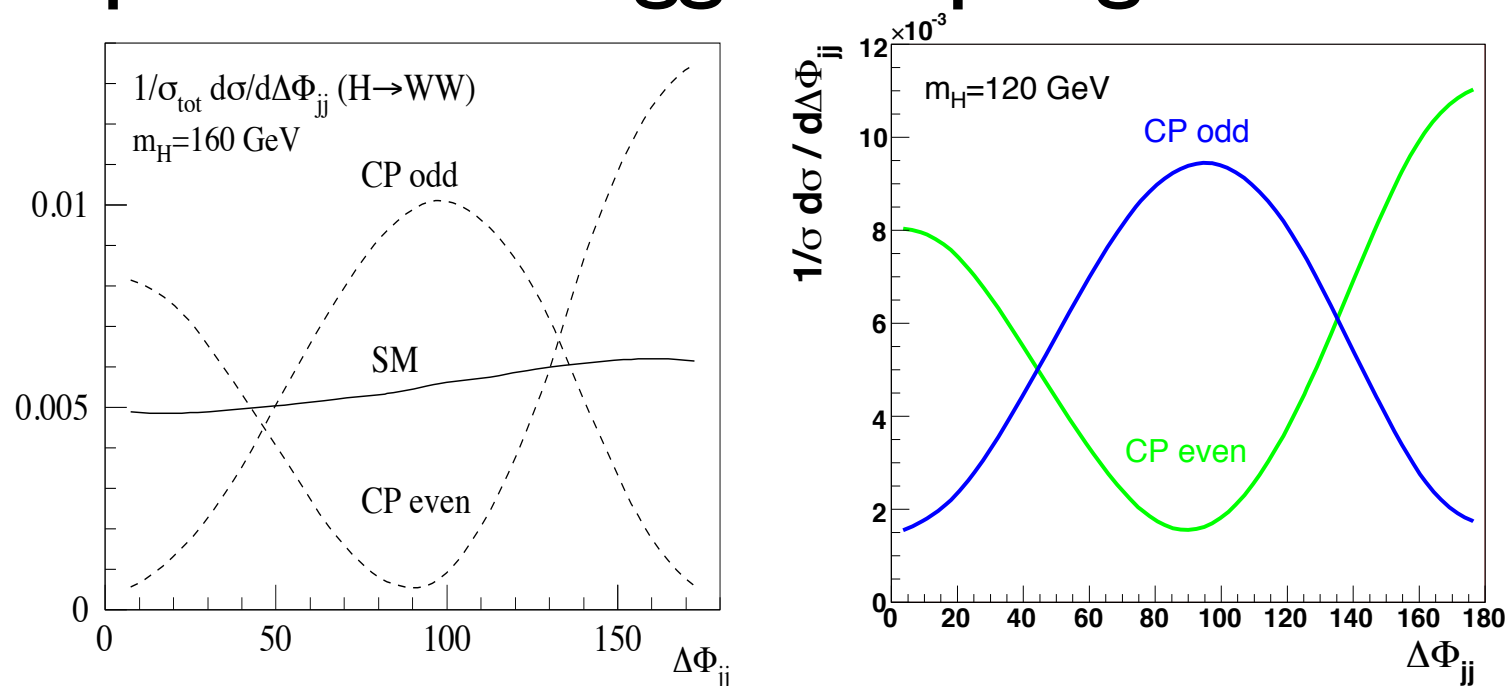
Inspiration from Higgs Search

- Proposal from Zeppenfeld *et. al.*
- Consider on-shell Higgs production from Vector Boson Fusion (VBF)
- Azimuthal angular dependence comes from gauge boson helicity
- Presence of various cos/sin modes depends on how these helicities can be combined.
- *i.e.* on the Lorentz structure of the matrix element for Higgs production



Inspiration from Higgs Search

- Searches for invisible Higgs decay
 - Look for azimuthal angular correlations in forward jets
 - Also shown that $d\sigma/d\Delta\phi$ sensitive to CP-properties of Higgs coupling

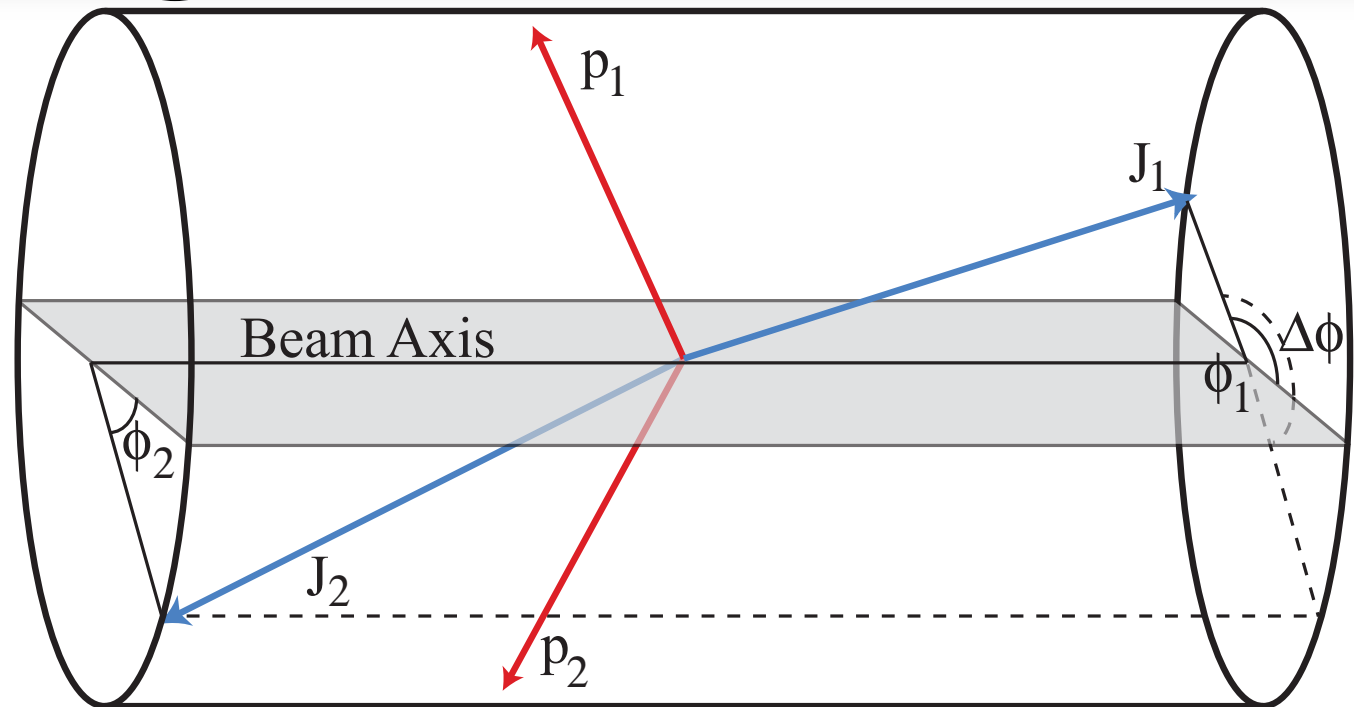


- Background has no $\cos 2\Delta\phi$ mode

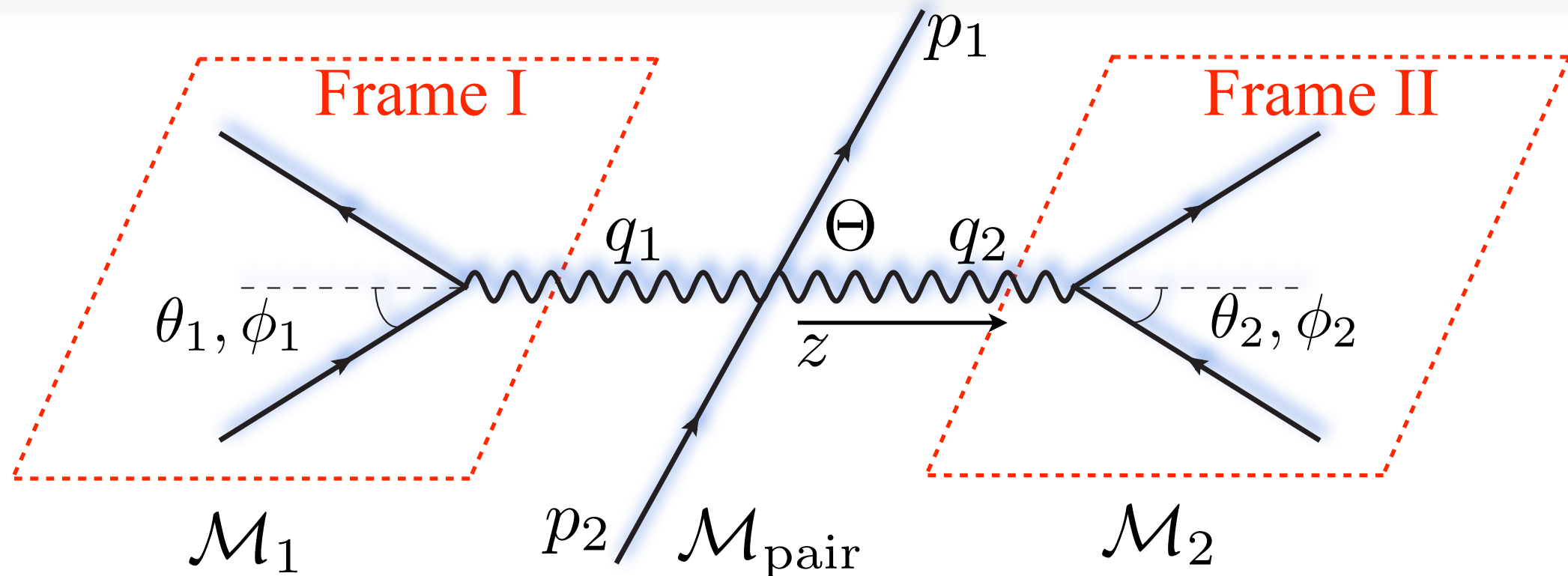
hep-ph/0105325, hep-ph/0605117, hep-ph/0703202

The Big Picture

- Since the spin measurement relies on the kinematics of jets J_1/J_2 , the only requirements on new physics (p_1/p_2) is that we can trigger on it (and identify the forward jets)
- For this introductory study, we assume both these problems can be ignored
 - Clearly, we are not experimentalists.



VBF Kinematics



- With these choices, ϕ dependence made clear:

$$\epsilon_{1/2} \propto e^{+i\phi_{1/2}}, e^{-i\phi_{1/2}}, e^{0 \times i\phi_{1/2}}$$

- I've drawn quark initial states only, but anti-quark and gluon contribute as well.

Azimuthal Angular Dependence

- Can expand out dependence on ϕ_1, ϕ_2 :

$$|\mathcal{M}|^2 = \left| \sum_{h_1, h_2 = \pm, 0} \mathcal{M}_1 \mathcal{M}_2 \mathcal{M}_{\text{pair}} e^{i(h_1 \phi_1 + h_2 \phi_2)} \right|^2$$

$$|\mathcal{M}|^2 \propto A_0 + A_1 \cos \Delta\phi + A_2 \cos 2\Delta\phi$$

After integrating over $\phi_1 + \phi_2$

A_2

- The coefficient A_1 gets a contribution from cuts.
- Look at the coefficient of $\cos 2\Delta\phi$ instead.

- From

$$|\mathcal{M}|^2 = \left| \sum_{h_1, h_2 = \pm, 0} \mathcal{M}_1 \mathcal{M}_2 \mathcal{M}_{\text{pair}} e^{i(h_1 \phi_1 + h_2 \phi_2)} \right|^2$$

we're interested in

$$A_2 = (\mathcal{PS}) \sum_{\substack{h_1, h'_1, h_2, h'_2 \\ |h_i - h'_i| = 2}} (\mathcal{M}_1(h_1) \mathcal{M}_2(h_2) \mathcal{M}_{\text{pair}}(h_1, h_2)) \times \\ (\mathcal{M}_1(h'_1) \mathcal{M}_2(h'_2) \mathcal{M}_{\text{pair}}(h'_1, h'_2))^*$$

Scalar Case

- Factoring out production matrix elements:

$$A_2 = (\mathcal{PS}) \mathcal{M}_1(+1) \mathcal{M}(-1)^* \mathcal{M}_2(-1) \mathcal{M}_2(+1)^* \\ [\mathcal{M}_{\text{pair}}(+1, -1) \mathcal{M}_{\text{pair}}(-1, +1)^* + (+1 \leftrightarrow -1)]$$

- In an abelian theory, easy to write down the matrix elements for transverse polarizations:

$$\mathcal{M}_{\text{scalar}} \propto (\epsilon_1 \cdot \epsilon_2) - 4 \left[\frac{(p_1 \cdot \epsilon_1)(p_1 - q_1) \cdot \epsilon_2}{q_1^2 - 2p_1 \cdot q_1} + \frac{(p_1 \cdot \epsilon_2)(p_1 - q_2) \cdot \epsilon_1}{q_2^2 - 2p_1 \cdot q_2} \right]$$

- Invariant under $\epsilon_{1/2}^+ \leftrightarrow \epsilon_{1/2}^-$
 - (Also true in non-abelian calculation)

$$A_2 \propto \mathcal{M}(+1, -1) \mathcal{M}(-1, +1)^* > 0$$

Spinor Case

- Straightforward for on-shell abelian example:

$$\mathcal{M}(\pm 1, \mp 1) = i\bar{u} \left[\frac{\not{\epsilon}_1^\pm (\not{p}_1 - \not{q}_1 + M) \not{\epsilon}_2^\mp}{q_1^2 - 2p_1 \cdot q_1} + \frac{\not{\epsilon}_2^\mp (\not{p}_1 - \not{q}_2 + M) \not{\epsilon}_1^\pm}{q_2^2 - 2p_1 \cdot q_2} \right] v$$

$$A_2 \propto -\frac{64m^2}{s} \left(1 + \frac{4m^2}{s\sqrt{1 - 4m^2/s}} \tanh^{-1} \sqrt{1 - 4m^2/s} \right) < 0$$

- Non-abelian example more subtle.
- Can divide $\mathcal{M}(+1, -1)\mathcal{M}(-1, +1)^*$ into symmetric ($d_{abc}d_{abd}T^cT^d$) and antisymmetric ($f^{abc}f^{abd}T^cT^d$) parts
 - Symmetric part reproduces abelian $A_2 < 0$
 - Asymmetric part naively gives $A_2 > 0$

Spinor Case

- Naive calculation ignores phase space cuts experimentally necessary to isolate VBF diagrams
 - (also not gauge invariant, as we calculate only VBF diagrams, not the full 1000+ possible)
- After cuts, isolating events with fusing gluons that are space and spin symmetric. Thus color-antisymmetric states don't contribute.
- Simulation through Calchep and MadGraph confirm that, for spinors,

$$A_2 < 0$$

Simulation Results

- Use MadGraph/MadEvent for background-free simulation:

$$pp \rightarrow 2(R - \text{hadrons}) + jj$$

- 500 GeV R-hadrons (excluded by Lepton/Photon)
- Apply VBF-isolating cuts:

$$\eta_{j_1} \cdot \eta_{j_2} < 0, \quad |\eta_j| \leq 5, \quad |\eta_{j_1} - \eta_{j_2}| \geq 4.2$$

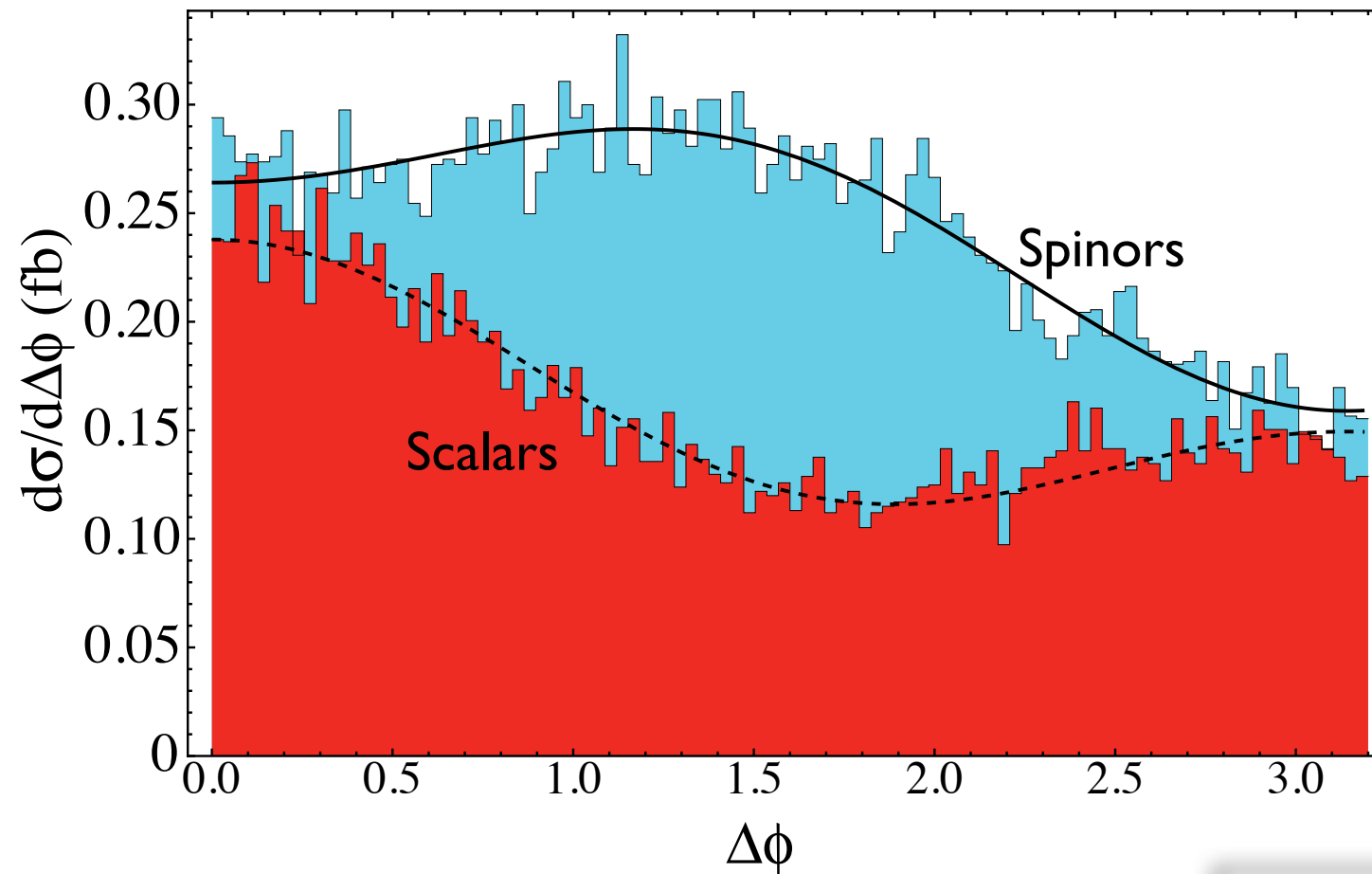
$$p_{T,j_1} \geq 30 \text{ GeV}, \quad p_{T,j} \geq 20 \text{ GeV}, \quad M_{jj} \geq 500 \text{ GeV}$$

$$|\eta_{R-\text{hadron}}| < 2.1, \quad p_{T,R-\text{hadron}} > 50 \text{ GeV}$$

- Total cross section ($\sqrt{s} = 10 \text{ TeV}, m = 500 \text{ GeV}$):

$$\sigma_{\text{spinor}} = 33 \text{ fb}, \quad \sigma_{\text{scalar}} = 21 \text{ fb}$$

Results



$$\frac{d\sigma_{\text{scalar}}}{d\Delta\phi} = 0.16 + 0.044 \cos \Delta\phi + 0.035 \cos 2\Delta\phi \quad (\text{fb})$$

$$(A_2/A_0)_{\text{scalar}} = 0.22$$

$$\frac{d\sigma_{\text{spinor}}}{d\Delta\phi} = 0.24 + 0.053 \cos \Delta\phi - 0.033 \cos 2\Delta\phi \quad (\text{fb})$$

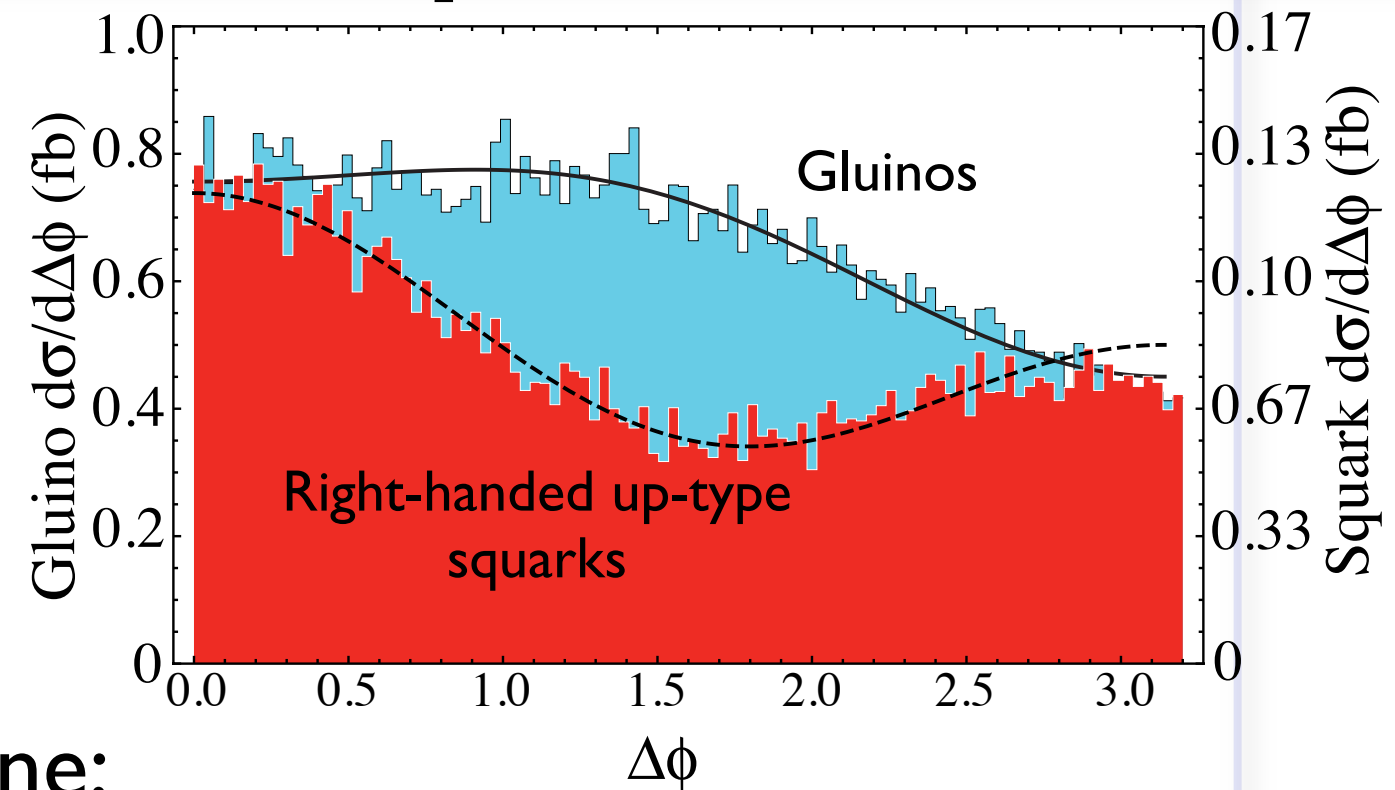
$$(A_2/A_0)_{\text{spinor}} = -0.14$$

SUSY Applications

- We picked a “background free” model
 - No central jets that can be confused with the forward jets that constitute our observables
- Obvious next step: SUSY gluino pairs/squark pairs
- What we have done:
 - Background-free, trigger/tagging free MadGraph study (*i.e.* is there a signal?)

Gluinos and Squarks

- Demonstrates that a signal is present, and that the majorana nature of the gluinos isn't a problem.
- What needs to be done:
 - Background (naive expectation: flat in $\cos 2\Delta\phi$)
 - Trigger analysis, cut optimization
 - Jet ID, including decays of gluinos/squarks



Problems with Pythia

- Background analysis (including decays) requires simulation with Pythia, as our signal relies on forward jets
- However, Pythia-generated jets do not include helicity information
 - MadGraph does, but the overall cross-section is wrong (no matching)
- Therefore, we claim that simulation of forward jets does not include correlations which contain useful physics information!

Future Work

- Understand how to correctly integrate MadGraph and Pythia results in the forward region
 - Also a useful test of the effect on cuts on $\Delta\phi$ distributions
- Optimize cuts to for signal & cross section
- Test at LHC using $t\bar{t}$ production?
- Examine vector pair production
- Look for physics information in $\phi_1 + \phi_2$ coefficients

Conclusions

- Correlations in forward jets originating in VBF events contain useful information about spin
- Allows a “model-independent” measurement
- This work reveals a kinematic region where existing simulation tools are inadequate
- Important information can exist in angular correlations, but not all simulators include these.

